

## Magnetic moment measurement of isomeric state in $^{75}\text{Cu}$

Y. Ichikawa,<sup>\*1</sup> A. Takamine,<sup>\*1</sup> H. Nishibata,<sup>\*2</sup> K. Imamura,<sup>\*1,\*3</sup> T. Fujita,<sup>\*1,\*2</sup> T. Sato,<sup>\*4</sup> S. Momiyama,<sup>\*5</sup> Y. Shimizu,<sup>\*1</sup> D. S. Ahn,<sup>\*1</sup> K. Asahi,<sup>\*4</sup> H. Baba,<sup>\*1</sup> D. L. Balabanski,<sup>\*1,\*6</sup> F. Boulay,<sup>\*1,\*7,\*8</sup> J. M. Daugas,<sup>\*1,\*8</sup> T. Egami,<sup>\*9</sup> N. Fukuda,<sup>\*1</sup> C. Funayama,<sup>\*4</sup> T. Furukawa,<sup>\*10</sup> G. Georgiev,<sup>\*11</sup> A. Gladkov,<sup>\*1,\*12</sup> N. Inabe,<sup>\*1</sup> Y. Ishibashi,<sup>\*1,\*13</sup> Y. Kobayashi,<sup>\*14</sup> S. Kojima,<sup>\*4</sup> A. Kusoglu,<sup>\*11,\*15</sup> T. Kawaguchi,<sup>\*9</sup> T. Kawamura,<sup>\*2</sup> I. Mukul,<sup>\*16</sup> M. Niikura,<sup>\*5</sup> T. Nishizaka,<sup>\*9</sup> A. Odahara,<sup>\*2</sup> Y. Ohtomo,<sup>\*1,\*4</sup> D. Ralet,<sup>\*11</sup> T. Shimoda,<sup>\*2</sup> G. S. Simpson,<sup>\*17</sup> T. Sumikama,<sup>\*1</sup> H. Suzuki,<sup>\*1</sup> H. Takeda,<sup>\*1</sup> L. C. Tao,<sup>\*1,\*18</sup> Y. Togano,<sup>\*4</sup> D. Tominaga,<sup>\*9</sup> H. Ueno,<sup>\*1</sup> H. Yamazaki,<sup>\*1</sup> and X. F. Yang<sup>\*19</sup>

The  $^{75}\text{Cu}$  nucleus has attracted much attention because the ground-state spin parity changes from  $3/2^-$  to  $5/2^-$  as a result of the migration of the  $5/2^-$  levels along the Cu isotopic chain<sup>2)</sup>. The  $^{75}\text{Cu}$  nucleus has two isomeric states<sup>3)</sup> at 62-keV and 66-keV levels directly decaying to the ground state<sup>4)</sup>, one of which is expected to have a spin parity of  $3/2^-$  inherited from the ground state of  $^{73}\text{Cu}$ . In order to investigate the wave function of the  $3/2^-$  state and to compare it with the  $5/2^-$  ground state<sup>1)</sup>, the magnetic moment of the isomeric state of  $^{75}\text{Cu}$  was measured.

The experiment was carried out at the BigRIPS at the RIBF. The two-step fragmentation scheme with momentum-dispersion matching<sup>5)</sup> was employed to produce highly spin-aligned  $^{75}\text{Cu}$ . In the reaction at F0,  $^{76}\text{Zn}$  was produced by a fission reaction of a 345-MeV/nucleon  $^{238}\text{U}$  beam on a  $^9\text{Be}$  target with a thickness of 1.29 g/cm<sup>2</sup>. The secondary  $^{76}\text{Zn}$  beam was introduced to a second target of wedge-shaped aluminum with a mean thickness of 0.81 g/cm<sup>2</sup>, placed at the momentum-dispersive focal plane F5. The  $^{75}\text{Cu}$  nuclei including those in isomeric state  $^{75\text{m}}\text{Cu}$  were produced through one-proton removal from  $^{76}\text{Zn}$ . The  $^{75}\text{Cu}$  beam was subsequently transported to F7 under the condition that the momentum dispersion between F5 and F7 was matched to that between F3 and F5.

The  $g$ -factor of  $^{75\text{m}}\text{Cu}$  was determined by means of the time-differential perturbed angular distribution (TDPAD) methods. The TDPAD apparatus, placed at F8, consists of a dipole magnet, a Cu crystal stop-

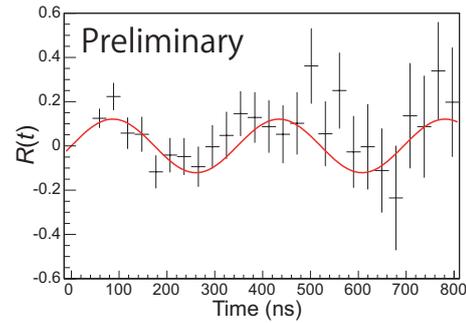


Fig. 1. Preliminary result of  $R(t)$  ratio for the 66-keV  $\gamma$  ray. The solid line represents the sine function after fitting to the the experimental  $R(t)$  ratio.

per, Ge detectors, and a plastic scintillator. The dipole magnet provided a static magnetic field of  $B_0 = 0.200$  T.  $^{75\text{m}}\text{Cu}$  was implanted into the Cu stopper, and  $\gamma$  rays were detected with four Ge detectors located perpendicular to  $B_0$  at a distance of 7.0 cm from the stopper and at angles of  $\pm 45$  and  $\pm 135$  degrees with respect to the beam axis. The plastic scintillator of 0.1 mm in thick was placed upstream of the stopper, the signal from which provided the time-zero trigger.

The  $R(t)$  ratio representing the change of anisotropy of  $\gamma$  rays synchronized with the spin precession was obtained according to

$$R(t) = \{N_{13}(t) - \epsilon N_{24}(t)\} / \{N_{13}(t) + \epsilon N_{24}(t)\}, \quad (1)$$

where  $N_{13}(t)$  and  $N_{24}(t)$  are the sums of the photopeak count rates at the two pairs of Ge detectors placed diagonally to each other, and  $\epsilon$  denotes a correction factor for the difference in the detection efficiency. In this experiment we observed an oscillation pattern only for the 66-keV  $\gamma$  ray with over  $5\sigma$  significance, as shown in Fig. 1. The magnitude of spin alignment was found to be larger than 50%. The detailed analysis and the deduction of the  $g$ -factor is in progress.

### References

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<sup>\*1</sup> RIKEN Nishina Center  
<sup>\*2</sup> Department of Physics, Osaka University  
<sup>\*3</sup> Department of Physics, Meiji University  
<sup>\*4</sup> Department of Physics, Tokyo Institute of Technology  
<sup>\*5</sup> Department of Physics, University of Tokyo  
<sup>\*6</sup> ELI-NP, IFIN-HH  
<sup>\*7</sup> GANIL, CEA/DSM-CNRS/IN2P3  
<sup>\*8</sup> CEA, DAM, DIF  
<sup>\*9</sup> Department of Advanced Sciences, Hosei University  
<sup>\*10</sup> Department of Physics, Tokyo Metropolitan University  
<sup>\*11</sup> CSNSM, CNRS/IN2P3, Université Pris-Sud  
<sup>\*12</sup> Department of Physics, Kyungpook National University  
<sup>\*13</sup> Department of Physics, University of Tsukuba  
<sup>\*14</sup> Department of Informatics and Engineering, University of Electro-Communications  
<sup>\*15</sup> Department of Physics, Istanbul University  
<sup>\*16</sup> Department of Particle Physics, Weizmann Institute  
<sup>\*17</sup> LPSC, Université Joseph Fourier Grenoble 1, CNRS/IN2P3  
<sup>\*18</sup> School of Physics, Peking University  
<sup>\*19</sup> Instituut voor Kern-en Srralingsfysica, K. U. Leuven