

# Atomic-number identification of heavy RI beams using the energy loss in a Xe-based gas

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Radioactive isotope (RI) beams produced at RIBF are tagged event-by-event with the atomic number  $Z$  and mass-to-charge ratio  $A/Q$  determined using beam-line detectors. For heavy RI beams, particle identification (PID) becomes difficult owing to the change in  $Q$  inside the beam-line detectors because  $Z$  is determined from the energy loss depending on  $Q^2$ . Blobs in the PID plot were clearly visible for the  $^{208}\text{Rn}$  case<sup>1)</sup> but not for the  $^{220}\text{Th}$  case.<sup>2)</sup> The relative  $Z$  resolutions were 0.45% and 0.69% ( $1\sigma$ ) for the 185-MeV/nucleon  $^{208}\text{Rn}$  beam and 315-MeV/nucleon  $^{220}\text{Th}$  beam, respectively. The worse resolution was considered to be due to the energy-loss straggling caused by charge-state fluctuation in the gas of the ionization chamber. The difference between these two cases indicates the impact of the energy dependence of the charge-state fluctuation.

In this paper, the difference in the  $Z$  resolutions is discussed in regards to the energy dependence of the cross section of the change in  $Q$  in the gas using the GLOBAL code.<sup>3)</sup> Figure 1 shows the energy dependence of the partial mean free path length  $L$ , given a change in  $Z - Q$  from 1 to 2 or from 2 to 1. The energy dependence of  $L$  is mainly for the electron-pickup reaction. In the  $^{208}\text{Rn}$  case, the mean value of the equilibrium charge-state distribution  $\langle Q \rangle$  is 84.5 at 185 MeV/nucleon. Since  $\langle Q \rangle$  is a decimal, the charge state must change multiple times in the ionization chamber to make the effective  $Q$  in a single event closer to 84.5.  $L$  ( $Z - Q = 1 \rightarrow 2$ ) is roughly 1/2 of the effective length of the ionization chamber, as indicated by the dotted line in Fig. 1. Thus, even if  $Z - Q$  changes from 2 to 1, it could be back. In contrast,  $L$  ( $Z - Q = 1 \rightarrow 2$ ) is longer than the effective length at 315 MeV/nucleon.  $Z - Q$  might not change to 2 once it becomes 1. This is mainly the nature of Ar gas, which accounts for 90% of the P10 gas. Figure 1 also shows  $L$

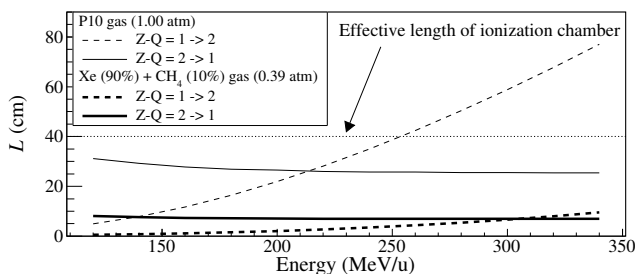


Fig. 1. Partial mean free path length of the change in the charge state in the ionization chamber as a function of the energy of the  $^{208}\text{Rn}$  beam.

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of a Xe-based mixed gas (Xe 90% + CH<sub>4</sub> 10%). The gas pressure was determined so that the energy deposit at 300 MeV/nucleon is same as that of the P10 gas. Even at 300 MeV/nucleon,  $L$  is shorter than that of P10 at 185 MeV/nucleon, indicating a better  $Z$  resolution.

The energy loss in the ionization chamber was simulated by using the energy-loss code ATIMA<sup>4)</sup> and GLOBAL. The fluctuation in  $Q$  was taken into account by the Monte Carlo method. The  $Z$  resolution of the 180-MeV/nucleon  $^{210}_{85}\text{At}$  beam was simulated to be 0.46%, which is consistent with the experimental value of 0.45%. Figure 2 shows the energy deposit of the 300-MeV/nucleon  $^{208}\text{Rn}$  and  $^{206}\text{Ac}$  beams into the effective region of the ionization chamber. The  $Z$  resolution is improved from 0.60% for the P10 gas to 0.39% for the Xe-based gas. This result is consistent with the discussion of Fig. 1.

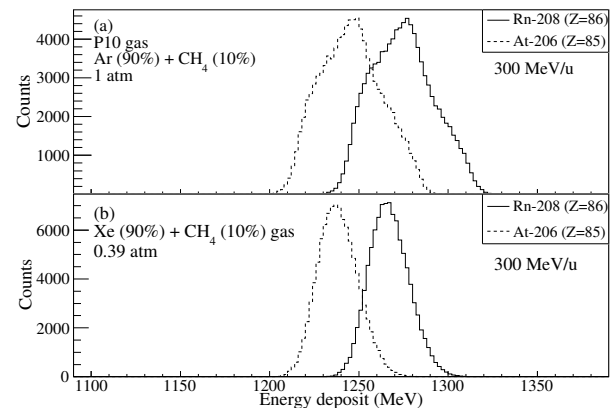


Fig. 2. Monte Carlo simulation of the energy deposit in the ionization chamber. Half of the energy-loss straggling in ATIMA was applied for the energy deposit.<sup>5)</sup>

An experimental study of PID using a Xe-based gas was already conducted, and a good  $Z$  resolution was obtained.<sup>6,7)</sup>

## References

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